

## Dielectric and Ferroelectric Properties of BaTiO<sub>3</sub> Nanofibers Prepared via Electrospinning



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BaTiO<sub>3</sub> nanofibers of about 400 nm in diameter were synthesized via electrospinning. The evolution of the morphology and phase composition of the BaTiO<sub>3</sub> nanofibers was studied by scanning electron microscopy and X-ray diffraction within the annealing temperature of 750–1050 °C. Higher annealing temperature led to rougher surface and better crystallization of the BaTiO<sub>3</sub> nanofibers. Below 1050 °C, the BaTiO<sub>3</sub> nanofibers maintained its large aspect ratios and could still be regarded as individual nanofiber. The dielectric permittivities of the BaTiO<sub>3</sub> nanofibers ( $\epsilon_r \sim 820$ ) were calculated with the MG equation by considering the porous bulk specimens as composites of BaTiO<sub>3</sub> nanofibers and air. The ferroelectric properties of the BaTiO<sub>3</sub> nanofibers were measured by using a ferroelectric analyzer coupled with an atomic force microscope. *P–E* loop measured for the BaTiO<sub>3</sub> nanofiber exhibits small hysteresis.

**KEY WORDS:** Dielectric; Ferroelectric; Barium titanate; Electrospinning

### 1. Introduction

BaTiO<sub>3</sub> is studied as a typical ferroelectric material with high dielectric permittivity, low dielectric loss and other excellent ferroelectric properties<sup>[1]</sup>. Owing to the outstanding properties, BaTiO<sub>3</sub> has been widely applied in capacitors, sensors, nonvolatile random access memory and many composites<sup>[2]</sup>. With their high intrinsic dielectric permittivities, the nanostructured BaTiO<sub>3</sub> particles, such as spheres<sup>[3]</sup> or fibers<sup>[4]</sup>, have also been employed as dielectric fillers in polymer composites to improve the dielectric behavior of the composites. The dielectric as well as geometric (size, shape etc) properties of the BaTiO<sub>3</sub> particles play a critical role in determining the dielectric behavior of the polymer composites<sup>[5,6]</sup>. There have also been efforts in making three-phase polymer composites filled with both BaTiO<sub>3</sub> particles and conductive fillers. The presence of BaTiO<sub>3</sub> nanoparticles prevents the conductive fillers from connecting into conducting paths, giving rise to lower dielectric loss. Our recent results show that a

kind of three-phase polymer composite filled with BaTiO<sub>3</sub> nanoparticles and graphene nanopellets could be polarized up to 150 kV/mm without much increased leakage current<sup>[7]</sup>. Yet, BaTiO<sub>3</sub> nanofibers are considered more promising over their spherical counterparts as dielectric fillers to improve the dielectric strength of their composites with polymers under high electric field. Due to their large aspect ratios and hence larger dipolar moment along their longitudinal axis, BaTiO<sub>3</sub> nanofibers are capable of increasing the dielectric permittivity of polymer composites<sup>[8]</sup>. Recent studies further indicate that surface treatments of BaTiO<sub>3</sub> nanofibers leads to stronger interfaces between the nanofibers and polymer matrix, giving rise to enhanced dielectric breakdown strength of polymer composites<sup>[9]</sup>. Enhanced dielectric strength of these BaTiO<sub>3</sub> nanofibers-filled polymer composites is attributed to the orientation of BaTiO<sub>3</sub> nanofibers in directions perpendicular to the external electric field. Both phase field modeling<sup>[6]</sup> and finite element simulation<sup>[10]</sup> indicate that the orientation of BaTiO<sub>3</sub> nanofibers is favorable for mitigating the electric field concentration in polymer matrix, which is caused by the large differences in dielectric permittivity between the BaTiO<sub>3</sub> ( $\epsilon_r \sim 1000$ ) and polymer matrix ( $\epsilon_r \sim 2-3$ ).

Electrospinning is one of the methods to fabricate one-dimensional materials. It is a convenient and versatile method to produce fibers, for instance, polymer, ceramic and composite fibers<sup>[11–14]</sup>. In this study, continuous BaTiO<sub>3</sub> nanofibers are prepared via electrospinning with sol–gel precursor.

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Dielectric and ferroelectric properties of BaTiO<sub>3</sub> ceramics have been intensively studied and well documented. The measurements are usually performed to dense ceramics sintered at relatively high temperature. There have been very limited reports on the dielectric behavior of BaTiO<sub>3</sub> nanofibers<sup>[15]</sup>. Although piezoelectric force microscopy (PFM) has been employed for the measurement of the piezoelectric response of the aligned BaTiO<sub>3</sub> nanofibers<sup>[16]</sup>, the measurement of their dielectric permittivities still requires a bulk specimen which is normally obtained through sintering at high temperature. The dielectric behavior obtained from a dense body thus could not be regarded as the intrinsic dielectric permittivities of the individual BaTiO<sub>3</sub> nanofibers. Here, in this contribution, we explored the evolution of morphology and phase composition of BaTiO<sub>3</sub> nanofibers with annealing temperature up to 1050 °C, the upper limit beyond which the BaTiO<sub>3</sub> nanofibers are sintered into bundles or cluster and can't be treated as individual nanofibers. Porous bulk specimens of the BaTiO<sub>3</sub> nanofibers are calcined at <1050 °C with their dielectric permittivities measured. The dielectric permittivities of the individual BaTiO<sub>3</sub> nanofibers are then calculated with a Maxwell–Garnett equation by considering the specimens as composites of BaTiO<sub>3</sub> nanofibers and air. The ferroelectric properties of the BaTiO<sub>3</sub> nanofibers are also investigated by using a ferroelectric test module coupled with an atomic force microscope (AFM).

## 2. Experimental

The chemicals were purchased from China National Medicines Corporation Ltd. if not otherwise specified. BaTiO<sub>3</sub> nanofibers were prepared via electrospinning combined with sol–gel method. Briefly, barium acetate, tetrabutyl titanate and acetylacetone were dissolved in acetic acid at a molar ratio of 1:1:2 and stirred to get homogeneous barium titanate precursor sol. Poly (vinyl pyrrolidone) (PVP,  $M = 1,300,000$ ) was added to the precursor for the control of sol viscosity. The sol was then transferred into a syringe and electrospun by applying an electric field of 1.5 kV/cm. Fibers consisting of PVP and barium titanate precursor via electrospinning were then calcined at 550 °C for 1.5 h to remove PVP. White powders were obtained. For direct comparison, BaTiO<sub>3</sub> particles were also prepared with the sol of barium titanate precursor and PVP by calcining at 550 °C for 1.5 h. Both the BaTiO<sub>3</sub> nanofibers and BaTiO<sub>3</sub> particles were annealed at 750 °C, 850 °C, 950 °C and 1050 °C for 6 h.

For the measurement of dielectric properties, the BaTiO<sub>3</sub> nanofibers were cold-pressed into pellets of 10 mm in diameter at 2 MPa. The resultant pellets were thermally treated at 750 °C, 850 °C, 950 °C and 1050 °C for 6 h, respectively. Ag paste was then painted on both sides of the pellet samples as electrodes for the measurement of dielectric properties. The apparent densities ( $\rho_a$ ) of the resultant sample pellets after each heat treatment were measured by Archimedes method. The theoretical densities ( $\rho_t$ ) of BaTiO<sub>3</sub> specimens were also determined from the XRD patterns. The relative density was then determined as  $\rho_a/\rho_t$ .

X-Ray diffraction (XRD, Bruker D8) and scanning electron microscopy (SEM, JEOL4500) were employed to investigate the crystal structure and morphology of BaTiO<sub>3</sub> fibers and particles. Dielectric properties were measured by employing an HP 4294A precision impedance analyzer (Agilent) in the frequency range of 100 Hz to 100 MHz at room temperature. For the measurement of the ferroelectric properties, the BaTiO<sub>3</sub> nanofibers were glued to a Pt/Si substrate with silver paste as bottom electrode and grounded. The probe of the atomic force microscope (AFM, PISA XE-100)

was connected to a ferroelectric test module (TF analyzer 2000, axiACCT) and used as top electrode. During the measurement, topography image of the BaTiO<sub>3</sub> nanofibers was firstly obtained by AFM image and was then used to navigate the probe to touch the top of the BaTiO<sub>3</sub> nanofiber. The polarization–electric field ( $P$ – $E$ ) loops were then measured by the ferroelectric test module at 20 Hz.

## 3. Results and Discussion

The effects of annealing temperature on the evolution of morphology of BaTiO<sub>3</sub> nanofibers are first characterized by SEM. Fig. 1(a)–(d) shows the SEM images of BaTiO<sub>3</sub> nanofibers calcined at different temperatures. As the temperature increases from 750 °C to 1050 °C, the diameter of the BaTiO<sub>3</sub> nanofibers gradually decreases from ~800 nm to ~400 nm, which is also accompanied by the change of surface morphology of the nanofibers from quite smooth at 750 °C to rather rough at 1050 °C. The morphology evolution of BaTiO<sub>3</sub> nanofibers with temperature are in line with the trend observed by others. Li *et al.* reported the same change in surface roughness of the BaTiO<sub>3</sub> nanofibers after annealing from 500 °C to 800 °C<sup>[17]</sup>. Within this temperature range, the BaTiO<sub>3</sub> nanofibers undergo the transformation from amorphous to crystalline. With the further increase of annealing temperature up to 1050 °C, the BaTiO<sub>3</sub> nanofibers begin to sinter, giving rise to even rougher surface and much less pores in the nanofibers. For comparison, the BaTiO<sub>3</sub> particles derived from the same precursor sol used in the electrospinning process are subjected to the same heat treatments with their SEM images shown in Fig. 1(e)–(h). The most distinctive change in morphology of the BaTiO<sub>3</sub> particles could be observed in the sample annealed at 1050 °C (Fig. 1(h)), where most of the particles have been sintered into large clusters.

In addition to the evolution of the morphology of BaTiO<sub>3</sub> nanofibers with annealing temperature, phase transformation from cubic to tetragonal is also revealed by X-Ray diffraction patterns, as shown in Fig. 2. On the XRD patterns of samples annealed at 750 °C and 850 °C, all diffraction peaks could be assigned to cubic phase (space group: Pm3m). Further increase in the annealing temperature to 950 °C and 1050 °C leads to the phase transformation from cubic to tetragonal (space group: P4mm), as evidenced by the split of (200) and (311) peaks of cubic phase (as better distinguished in Fig. 2(b) and (c)). The same phase transformation is observed in the BaTiO<sub>3</sub> particles after the same heat treatments, with the only difference being the higher tetragonality of the BaTiO<sub>3</sub> nanofibers than that of the BaTiO<sub>3</sub> particles. As calculated from the XRD patterns,  $c/a$  ratio of BaTiO<sub>3</sub> fibers are 1.0083 at 950 °C and 1.0084 at 1050 °C, as compared with  $c/a$  of ~1.0076 at 950 °C and 1.0081 at 1050 °C for BaTiO<sub>3</sub> particles. Also It could be observed from the XRD patterns that the intensities of the (002)/(200) peaks and (113)/(311) peaks are much higher in the BaTiO<sub>3</sub> nanofibers than those in the BaTiO<sub>3</sub> particles (as shown in Fig. 2(b)–(f)), which is indicative of better crystallization after the annealing. The better crystallization of BaTiO<sub>3</sub> nanofibers is further evidenced by their larger grain sizes as compared with the BaTiO<sub>3</sub> particles. Fig. 3 shows the variation of grain size determined from the XRD patterns with annealing temperature. As seen in both the BaTiO<sub>3</sub> nanofibers and BaTiO<sub>3</sub> particles, the grain size consistently increases with increasing annealing temperature. The grain size of the BaTiO<sub>3</sub> nanofibers is larger than that of the BaTiO<sub>3</sub> particles at all temperatures.

Dielectric behaviors of both the BaTiO<sub>3</sub> nanofibers and particles are measured from the porous bulk specimen. With the

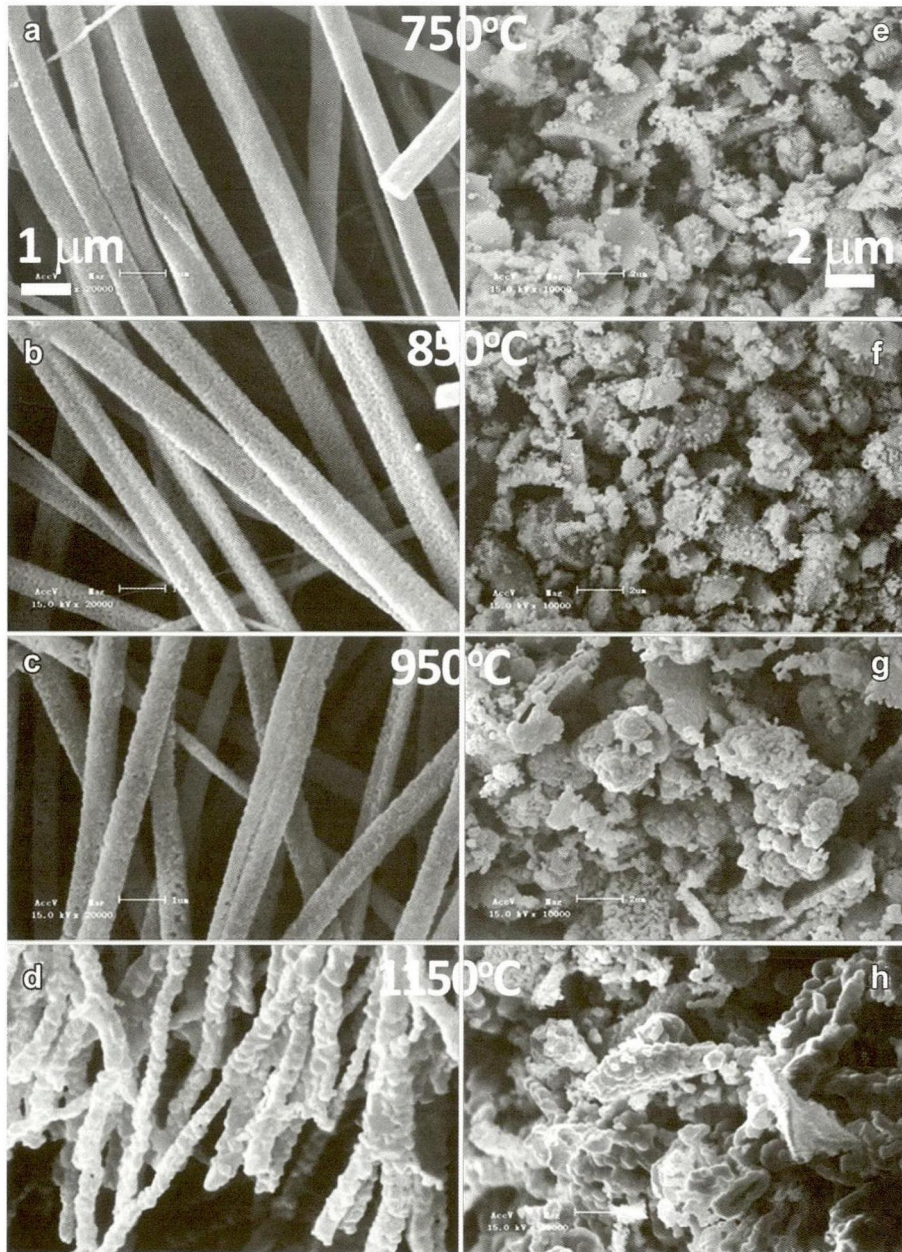


Fig. 1 SEM images of the BaTiO<sub>3</sub> nanofibers/particles calcined at 750 °C (a and e), 850 °C (b and f), 950 °C (c and g) and 1050 °C (d and h) for 6 h.

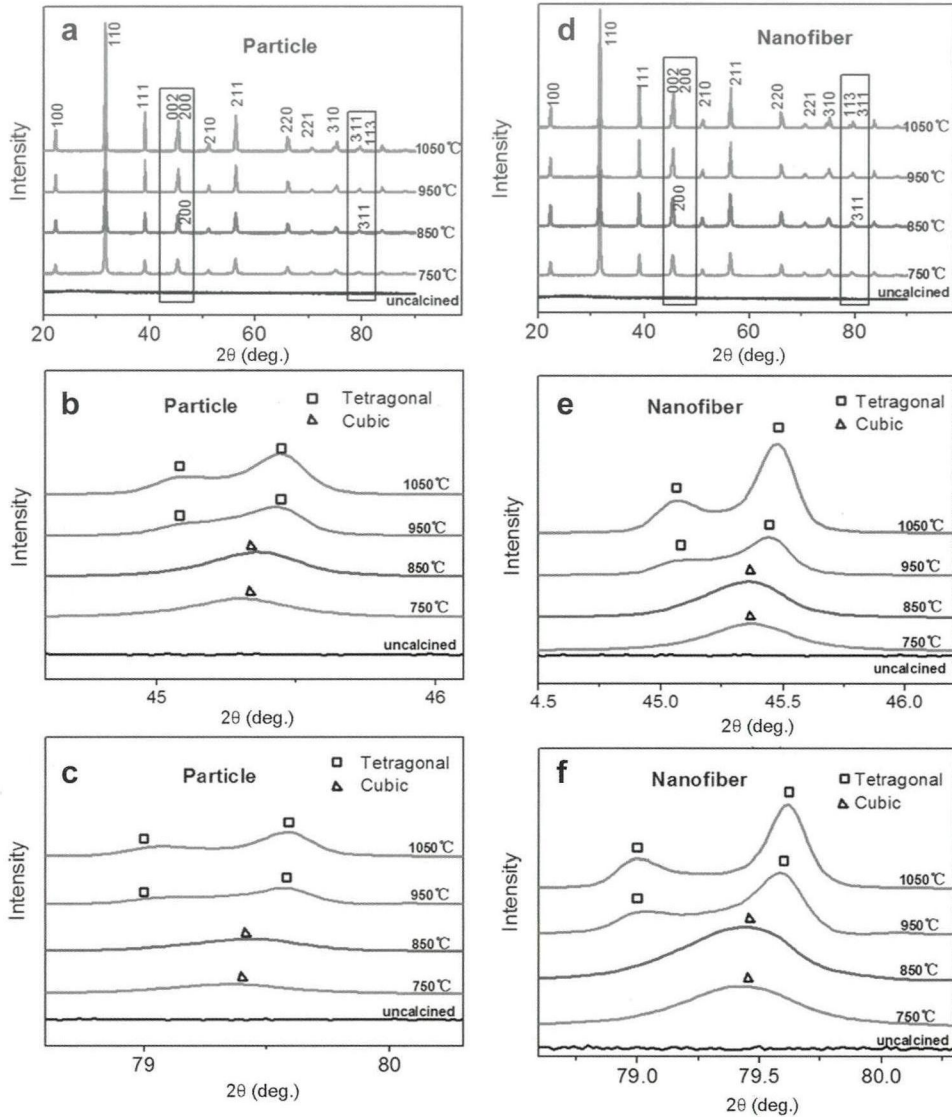
presence of pores in the specimen, the dielectric permittivities of these specimens are rather low, i.e., only  $\sim 200$  even for the specimen that has been annealed at 1050 °C. Two features can be readily observed from the frequency dependence of dielectric permittivities shown in Fig. 4(a). First, the dielectric permittivities remain stable within the frequency range of 100–10 MHz, suggesting no dielectric relaxation processes. Second, the dielectric permittivity increases with increasing annealing temperature for both the BaTiO<sub>3</sub> nanofibers and particles. In addition, the dielectric permittivities of the BaTiO<sub>3</sub> nanofibers are consistently higher than their particle counterparts at all temperatures. For the determination of dielectric permittivities of the BaTiO<sub>3</sub> nanofibers, the porous specimens are considered as composites of BaTiO<sub>3</sub> nanofibers and air. The dielectric

permittivity of the BaTiO<sub>3</sub> nanofibers can then be calculated by the Maxwell-Garnett equation as<sup>[18]</sup>:

$$\frac{\epsilon_r - \epsilon_i}{\epsilon_r + 2\epsilon_i} = f_a \cdot \frac{\epsilon_a - \epsilon_i}{\epsilon_a + 2\epsilon_i} \quad (1)$$

where  $\epsilon_r$  is the apparent dielectric permittivity of the specimen,  $\epsilon_i$  the intrinsic dielectric permittivity of the BaTiO<sub>3</sub> nanofibers,  $\epsilon_a$  the air dielectric permittivity of  $\sim 1$ .  $f_a = 1 - \rho_a/\rho_i$  is the volume fraction of air. The intrinsic dielectric permittivities calculated by the MG equation are plotted as a function of annealing temperatures in Fig. 4(b). In general, higher annealing temperature results in higher dielectric permittivities, both for the BaTiO<sub>3</sub> nanofibers and particles. Highest dielectric permittivity of



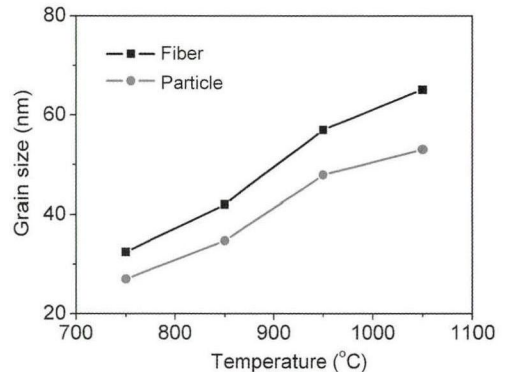


**Fig. 2** X-Ray diffraction patterns for the BaTiO<sub>3</sub> nanofibers (a) and the BaTiO<sub>3</sub> particles (b) after annealing at different temperatures. The phase transformation from cubic to tetragonal is better distinguished by the splitting of (200) peak into (200)/(002) peaks (c & d) and (311) peak into (311)/(113) peaks (e & f).

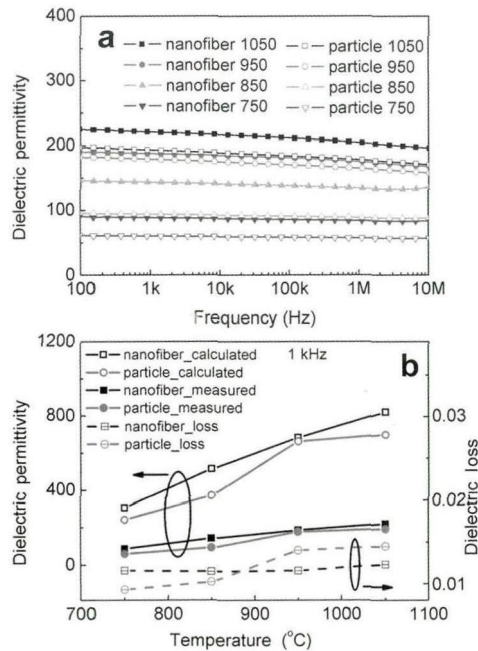
~ 820 is obtained for the BaTiO<sub>3</sub> nanofibers annealed at 1050 °C, which is comparable with the dielectric permittivities of the dense BaTiO<sub>3</sub> ceramic with an average grain size of ~ 20 nm<sup>[19]</sup>. Given the larger grain size of the BaTiO<sub>3</sub> nanofibers, the dielectric permittivity of the BaTiO<sub>3</sub> nanofibers, is still much smaller than the dielectric permittivity of the well sintered BaTiO<sub>3</sub> ceramics (usually >1000)<sup>[20]</sup>. Two factors have to be considered when discussing these results. The first one is that the BaTiO<sub>3</sub> nanofibers are treated at temperatures much lower than the sintering temperature. The interface between the BaTiO<sub>3</sub> nanofibers and air could have adverse effects on the dielectric behavior of the specimen. Moreover, given the rather high volume fraction of air in the porous specimen, the MG equation could underestimate the real dielectric permittivities of the BaTiO<sub>3</sub> nanofibers<sup>[18]</sup>. The dielectric permittivities thus determined should be considered as the lower limit for the intrinsic dielectric permittivities of the BaTiO<sub>3</sub> nanofibers.

Fig. 5 shows the polarization–electric field (*P*–*E*) loop for the BaTiO<sub>3</sub> nanofiber. The *P*–*E* loop exhibits very small hysteresis with almost linear dependence of *P*(*E*), more of the paraelectric behavior

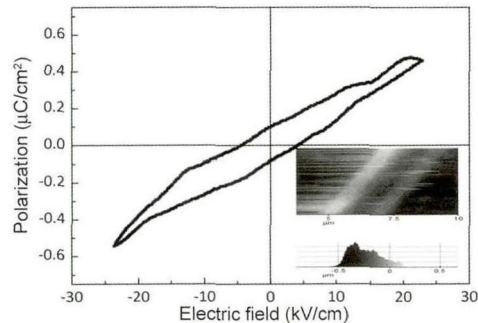
than ferroelectric. The strong suppression of macroscopic ferroelectric character for the BaTiO<sub>3</sub> nanofibers could be attributed to the following causes. First, the frozen domain structure under an external field owing to the effects exerted by the grain boundaries,



**Fig. 3** Variation of the grain sizes of BaTiO<sub>3</sub> nanofibers and particles with annealing temperatures.



**Fig. 4** (a) Dependences of dielectric permittivity on frequency from 100 Hz to 10 MHz of BaTiO<sub>3</sub> nanofibers and particles after annealing at different temperatures. (b) Dielectric permittivities calculated by the MG equation for BaTiO<sub>3</sub> nanofibers and particles after annealing at different temperatures at 1 kHz. The variation of dielectric loss with annealing temperatures is also superimposed.



**Fig. 5** Polarization–electric field loop for the BaTiO<sub>3</sub> nanofiber annealed at 1050 °C. The inset shows the topography image of the BaTiO<sub>3</sub> nanofiber measured.

such as the clamping of the domain walls and the hindrance of polarization switching, plays a key role in the suppression of ferroelectric hysteresis and switching<sup>[21]</sup>. Next, the depolarization field originated by the low permittivity nonferroelectric grain boundaries can cause a significant reduction of polarization<sup>[22]</sup>. It is also worth noting that the  $P$ – $E$  loop obtained in this manner only represents the polarization behavior in the direction of the short axle of the BaTiO<sub>3</sub> nanofibers. With the large differences in the depolarization field along the longitudinal direction and the lateral directions of the BaTiO<sub>3</sub> nanofibers, as calculated by Wang et al.<sup>[6]</sup>, the polarization behaviors of the BaTiO<sub>3</sub> nanofibers along the longitudinal direction might be much different.

#### 4. Conclusion

In summary, BaTiO<sub>3</sub> nanofibers of diameters around 400 nm are synthesized via electrospinning, and their dielectric and

ferroelectric properties are investigated. Below 1050 °C, BaTiO<sub>3</sub> nanofibers maintain its large aspect ratio and still can be regarded as individual fiber. BaTiO<sub>3</sub> nanofibers are better crystallized as compared with the BaTiO<sub>3</sub> particles after the same heat treatments. The dielectric permittivities of the BaTiO<sub>3</sub> nanofibers are calculated with the MG equation by considering the porous bulk specimens as composites of BaTiO<sub>3</sub> nanofibers and air. The  $P$ – $E$  loop measured for the BaTiO<sub>3</sub> nanofiber exhibits small hysteresis. Further investigation is still needed for the determination of the polarization behavior of the BaTiO<sub>3</sub> nanofibers along their longitudinal direction.

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#### REFERENCES

- [1] M.H. Frey, D.A. Payne, *Phys. Rev. B* 54 (1996) 3158–3168.
- [2] S. O'Brien, L. Brus, C.B. Murray, *J. Am. Chem. Soc.* 123 (2001) 12085–12806.
- [3] L. Li, A. Takahashi, J. Hao, R. Kikuchi, T. Hayakawa, T. Tsurumi, M. Kakimoto, *IEEE Trans. Comp. Packaging Technol.* 28 (2005) 754–759.
- [4] Y. Song, Y. Shen, H.Y. Liu, Y.H. Lin, M. Li, C.W. Nan, *J. Mater. Chem.* 22 (2012) 8063–8068.
- [5] Y.U. Wang, *Appl. Phys. Lett.* 96 (2010) 232901.
- [6] Y.U. Wang, D.Q. Tan, J. Krahn, *J. Appl. Phys.* 110 (2011) 044103.
- [7] Y. Shen, Y.H. Guan, Y.H. Hu, Y.H. Lin, C.W. Nan, *Appl. Phys. Lett.* 103 (2013) 072906.
- [8] H.A. Avila, L.A. Ramajo, M.S. Goes, M.M. Roberedo, M.S. Castro, R. Parra, *ACS Appl. Mater. Interfaces* 5 (2013) 505–510.
- [9] Y. Song, Y. Shen, H.Y. Liu, Y.H. Lin, M. Li, C.W. Nan, *J. Mater. Chem.* 22 (2012) 16491–16496.
- [10] Z.P. Wang, J.K. Nelson, H. Hillborg, S. Zhao, L.S. Schadler, *Comp. Sci. Technol.* 76 (2013) 29–36.
- [11] D. Li, J.T. McCann, Y.N. Xia, *J. Am. Chem. Soc.* 89 (2006) 1861–1869.
- [12] H. Wu, D.D. Lin, W. Pan, *Appl. Phys. Lett.* 89 (2006) 133125.
- [13] W.E. Teo, S. Ramakrishna, M. Kotaki, X.M. Mo, *Nanotechnology* 16 (2005) 918–924.
- [14] B. Sahoo, P.K. Panda, *Ceram. Int.* 38 (2013) 5189–5193.
- [15] W.S. Yun, J.J. Urban, Q. Gu, *Nano Lett.* 2 (2002) 447–450.
- [16] P. Sà, J. Barbosa, I. Bdikin, B. Almeida, A.G. Rolo, E.M. Gomes, M. Belsley, A.L. Kholkin, D. Isakov, *J. Phys. D: Appl. Phys.* 46 (2013) 105304.
- [17] H.P. Li, H. Wu, D.D. Lin, W. Pan, *J. Am. Ceram. Soc.* 92 (2009) 2162–2164.
- [18] C.W. Nan, *Prog. Mater. Sci.* 37 (1993) 1–116.
- [19] X.Y. Deng, X.H. Wang, H. Wen, L.L. Chen, L. Chen, L.T. Li, *Appl. Phys. Lett.* 88 (2006) 252905.
- [20] Z. Zhao, V. Buscaglia, M. Viviani, M.T. Buscaglia, L. Mitoseriu, A. Testino, M. Nygren, M. Johnsson, P. Nanni, *Phys. Rev. B* 70 (2004) 024107.
- [21] M.H. Frey, Z. Xu, P. Han, D.A. Payne, *Ferroelectrics* 337 (1998) 206–207.
- [22] M.T. Buscaglia, M. Viviani, V. Buscaglia, L. Mitoseriu, A. Testino, P. Nanni, Z. Zhao, M. Nygren, C. Harnagea, D. Piazza, C. Galassi, *Phys. Rev. B* 73 (2006) 064114.